

Fixation times in evolutionary games under weak selection

Philipp M. Altrock and Arne Traulsen

Max Planck Institute for Evolutionary Biology, August-Thienemann-Str. 2, 24306 Plön, Germany

altrock@evolbio.mpg.de, traulsen@evolbio.mpg.de

Abstract.

In evolutionary game dynamics, reproductive success increases with the performance in an evolutionary game. If strategy A performs better than strategy B , strategy A will spread in the population. Under stochastic dynamics, a single mutant will sooner or later take over the entire population or go extinct. We analyze the mean exit times (or average fixation times) associated with this process. We show analytically that these times depend on the payoff matrix of the game in an amazingly simple way under weak selection, i. e. strong stochasticity: The payoff difference $\Delta\pi$ is a linear function of the number of A individuals i , $\Delta\pi = u i + v$. The unconditional mean exit time depends only on the constant term v . Given that a single A mutant takes over the population, the corresponding conditional mean exit time depends only on the density dependent term u . We demonstrate this finding for two commonly applied microscopic evolutionary processes.

1. Introduction

Systems in which successful strategies spread by imitation or genetic reproduction can be described by evolutionary game theory. Such models are routinely analyzed in evolutionary biology, sociology, anthropology and economics. Recently, the application of methods from statistical physics to these systems has lead to many important insights [1, 2, 3, 4, 5].

Traditionally, the dynamics is described by the replicator equations, where the growth rate of a strategy is associated with its relative success compared with the population average [6, 7].

In the past years, research has focused on stochastic evolutionary game dynamics in finite populations [8, 9, 10, 11, 12, 13, 14, 15, 16, 17, 18, 19, 20, 21, 22]. In this context, a connection to the weak selection limit of population genetics has been established [8]. Weak selection means that the payoff differences based on different strategic behavior in interactions represent only a small correction to otherwise random dynamics, similar to high temperature expansions in physics. Weak selection is considered as a relevant limit in biology, as most evolutionary changes are driven by small fitness differences [23]. Moreover, it allows analytical approximations that are often impossible when selective differences in payoffs are large [8, 24, 25].

Most of the recent work that uses the weak selection approximation has been focusing on the probability that a certain strategy takes over. The time associated with this process has been calculated [26], but it received considerably less attention so far. Here, we present the weak selection corrections to the conditional and unconditional mean exit or fixation times in evolutionary 2×2 games with N players.

The conditional average time to fixation t_1^A is the expected time a single mutant needs to take over the population, given that such a takeover occurs at all. The unconditional average time of fixation t_1 is the expectation value for the time until the population is homogenous again after the arrival of a single mutant. This is regardless of whether the mutant type takes over the population or becomes extinct. Equivalently, the average fixation times for such one dimensional random walks can also be interpreted as mean first passage times or mean exit times [27, 28, 29].

Throughout this paper, we use the payoff matrix

$$\begin{array}{cc} & \begin{array}{cc} A & B \end{array} \\ \begin{array}{c} A \\ B \end{array} & \begin{pmatrix} a & b \\ c & d \end{pmatrix} \end{array}. \quad (1)$$

An A player interacting with another A receives a . If it interacts with B , it obtains b . Similarly, B receives c from A and d from other B 's. Thus, the average payoffs are

$$\pi_A(i) = \frac{i-1}{N-1} a + \frac{N-i}{N-1} b \quad (2)$$

$$\pi_B(i) = \frac{i}{N-1} c + \frac{N-i-1}{N-1} d. \quad (3)$$

A quantity that is of particular interest is the difference between the average payoffs,

$$\Delta\pi(i) = \pi_A(i) - \pi_B(i) = u i + v, \quad (4)$$

where

$$u = \frac{a + d - (b + c)}{N - 1}, \quad (5)$$

$$v = \frac{N(b - d) - (a - d)}{N - 1}. \quad (6)$$

We show that under weak selection, the conditional time (t_1^A) during which a single mutant takes over the whole population depends only on u (and, of course, on the population size). The unconditional time (t_1) during which the mutant either takes over the population or reaches extinction depends only on v (and the population size). See Figure 1 for an illustration of the relevant quantities.

Our manuscript is organized as follows: In Section 2, we introduce a particular evolutionary process for our analysis. Although our results are valid for a broader class of processes, we only present the full calculation for this evolutionary process. In Section 3, we recall the general form of fixation probabilities and times. We discuss neutral selection in Section 4 as a prerequisite to the weak selection expansion, which we explore in Section 5. In Section 6 we address the frequency dependent Moran process to underline the generality of our findings. The consequences of our analytical results are discussed in Section 7.

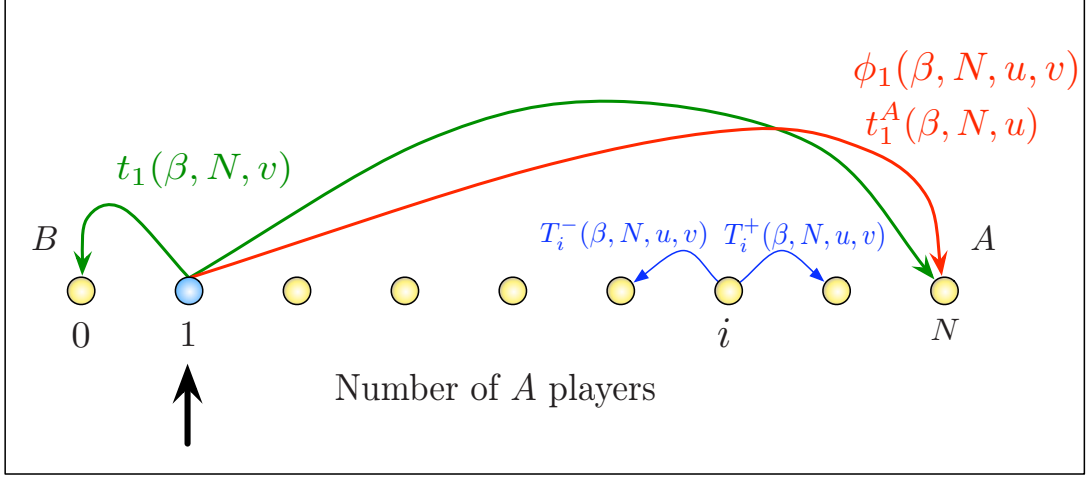


Figure 1: Illustration of the most relevant quantities. We are interested in the evolutionary fate of a single A player. All quantities depend on the intensity of selection β and the population size N . The payoff difference between A and B players is given by $\Delta\pi = u i + v$, with i as the number of A players. Both the transition probabilities T_i^+ and T_i^- and the probability that a single A player takes over the population ϕ_1 depend on u and v . But for weak selection, $\beta \ll 1$, the conditional time t_1^A during which a single A player takes over a population of B players only depends on u , whereas the unconditional time t_1 until either A or B has taken over the population only depends on v .

2. Fermi process

In a finite population of size N with two possible strategies A and B , the state of the system is characterized by the number of type A individuals i . In general, the dynamics is stochastic. In each time step, a randomly chosen individual evaluates its success. It compares this payoff with a second, randomly chosen individual. If this second individual has a higher payoff, the first one switches strategies with probability $p > \frac{1}{2}$. Otherwise, it switches with $p < \frac{1}{2}$. We assume that the switching probability is given by the Fermi distribution. Its shape is controlled by the intensity of selection β , which can be interpreted as an inverse temperature,

$$p_i^\pm = \frac{1}{1 + e^{\mp\beta(\pi_A(i) - \pi_B(i))}} = \frac{1}{1 + e^{\mp\beta\Delta\pi(i)}}. \quad (7)$$

In previous work [30, 31, 32], there is a different strategy update procedure. The first individual switches to the second's strategy with probability p_i^\pm . The second

individual can also switch to the first individual's strategy with probability $1 - p_i^\pm$. This yields a factor 2 in the transition probabilities (and, as we will become clear later, a factor $\frac{1}{2}$ in the fixation times). This process also has a proper strong selection limit, i.e. it is possible to examine $\beta \rightarrow \infty$. In this latter case we have $p_i^\pm \rightarrow \Theta(\Delta\pi(i))$, where $\Theta(x)$ is the step function.

The population size is constant in time, in each time step the state of the system can at most change by one, i.e. from i to $i-1$ or to $i+1$. The transition probabilities T_i^\pm to move from i to $i \pm 1$ are

$$T_i^\pm = \frac{i}{N} \frac{N-i}{N} p_i^\pm. \quad (8)$$

The probability to stay in the current state is $1 - T_i^+ - T_i^-$. An important measure of where the system is more likely to move is their ratio,

$$\gamma_i = \frac{T_i^-}{T_i^+} = e^{-\beta \Delta\pi(i)}. \quad (9)$$

This is a quantity that describes the tendency to move from the state i to $i \mp 1$, depending on whether $\gamma_i \gtrless 1$. Of course, $T_i^+ > 0$ is required, which follows from $\beta < \infty$. The T_i^\pm and thus the γ_i are invariant under adding a value to each of the payoffs given in (1), whereas multiplying the payoff matrix with a factor λ results in a change in the intensity of selection $\tilde{\beta} = \beta \lambda$.

Let us now focus on weak selection, $\beta \ll 1$. In this case we have

$$p_i^\pm \approx \frac{1}{2} \pm \frac{\beta}{4} \Delta\pi(i). \quad (10)$$

Weak selection corresponds to high temperature in Fermi statistics. A Taylor expansion of the γ_i up to first order in β yields $\gamma_i \approx 1 - \beta \Delta\pi(i)$. In this case, the probability to move from i to $i+1$ is very similar to the probability to move from i to $i-1$. Weak selection links the Fermi process to a variety of birth death processes, cf. [8, 33].

3. Fixation probabilities and fixation times

From equation (8) it follows that the two pure states all A or all B are absorbing, $T_0^\pm = T_N^\pm = 0$. In a finite population, we can calculate the probability ϕ_i that the system will fixate to the pure state all A , starting with the mixed state i . Obviously, we have $\phi_0 = 0$ and $\phi_N = 1$. For $0 < i < N$, there is a balance equation for the

fixation probabilities, $\phi_i = T_i^- \phi_{i-1} + (1 - T_i^+ - T_i^-) \phi_i + T_i^+ \phi_{i+1}$. This recursion leads to an expression for the fixation probabilities in terms of the γ_i [34, 35, 36],

$$\phi_i = \frac{1 + \sum_{k=1}^{i-1} \prod_{l=1}^k \gamma_l}{1 + \sum_{k=1}^{N-1} \prod_{l=1}^k \gamma_l}, \quad (11)$$

which is valid for any birth death process.

For the Fermi process, the exact equation (9) simplifies matters in an elegant way because the products in equation (11) can be solved,

$$\prod_{l=1}^k \gamma_l = \exp \left\{ -\beta \sum_{l=1}^k \Delta\pi(l) \right\} = \exp \left\{ -\beta \left[k^2 \frac{u}{2} + k \left(\frac{u}{2} + v \right) \right] \right\}. \quad (12)$$

Hence, equation (11) simplifies to

$$\phi_i = \frac{1 + \sum_{k=1}^{i-1} \exp \left\{ -\beta \left[k^2 \frac{u}{2} + k \left(\frac{u}{2} + v \right) \right] \right\}}{1 + \sum_{k=1}^{N-1} \exp \left\{ -\beta \left[k^2 \frac{u}{2} + k \left(\frac{u}{2} + v \right) \right] \right\}}. \quad (13)$$

For large N , the sums in equation (13) can be approximated by integrals, which yields a closed expression for the probabilities ϕ_i [33, 37].

General expressions for the unconditional and conditional mean exit times or average times of fixation, t_1 and t_1^A , are well known, especially for simple, translational invariant random walks [26, 27, 38]. A complete derivation for the average times of fixation in finite systems without translational invariance can be found in [26, 35, 39].

In the following, we will focus on the fixation of a single A mutant in a population of B . Accordingly, the unconditional and conditional fixation times read

$$t_1 = \phi_1 \sum_{k=1}^{N-1} \sum_{l=1}^k \frac{1}{T_l^+} \prod_{m=l+1}^k \gamma_m, \quad (14)$$

and

$$t_1^A = \sum_{k=1}^{N-1} \sum_{l=1}^k \frac{\phi_l}{T_l^+} \prod_{m=l+1}^k \gamma_m, \quad (15)$$

respectively. Time is measured in elementary time steps here. Thus, in each time step one reproductive event occurs. In biological contexts, it is often more convenient to measure time in generations, such that each individual reproduces once per generation on average. Time in generations is obtained by dividing the number of time steps by the population size N . It is well known that the variance of the

exit times under weak selection can be large [39], which has important biomedical implications [40]. Nonetheless, here we concentrate on the expectation values and do not address the distribution of the exit times.

4. Neutral selection

An important reference case is neutral selection, which results from vanishing selection intensity $\beta = 0$ [41]. Neutral selection is a very general limit, which is typically not affected by the details of the evolutionary process. For neutral selection we have $\gamma_i = 1$, that is $T_i^+ = T_i^-$ in any state i . However, we still have $T_i^\pm \neq T_j^\pm$ for $i \neq j$, although the system is symmetric, $T_i^\pm = T_{N-i}^\pm$. This is a difference to the simple random walk in one dimension, which is invariant with respect to translation [29].

For the Fermi process, the neutral transition probabilities are

$$T_i^\pm \Big|_{\beta=0} = \frac{1}{2} \frac{i}{N} \frac{N-i}{N}. \quad (16)$$

We have $T_i^+ = T_i^-$, which leads to $\gamma_i = 1$. From equation (11), it is thus clear that the probability of fixation to A is given by the initial abundance of A ,

$$\phi_i \Big|_{\beta=0} = \frac{i}{N}. \quad (17)$$

For the neutral unconditional time of fixation t_1 we get

$$t_1 \Big|_{\beta=0} = \frac{1}{N} \sum_{k=1}^{N-1} \sum_{l=1}^k \frac{2 N^2}{l(N-l)} = 2 N H_{N-1}. \quad (18)$$

Details for this calculation can be found in Appendix A. We introduced the shorthand notation for the harmonic numbers $H_{N-1} = \sum_{l=1}^{N-1} \frac{1}{l}$, which diverge logarithmically with N . In the same way we can solve

$$t_1^A \Big|_{\beta=0} = \sum_{k=1}^{N-1} \sum_{l=1}^k \frac{l}{N} \frac{2 N^2}{l(N-l)} = 2 N(N-1). \quad (19)$$

For neutral selection, the conditional average time of fixation of a single mutant diverges quadratically with the system size.

5. Weak Selection

In this section we will calculate the linear corrections of the mean exit times or fixation times t_1 , and t_1^A under weak selection, $\beta \ll 1$. Of course, all weak selection approximations are valid only if the term linear in β is small compared to the constant term.

The fixation probabilities for small β are

$$\phi_i \approx \frac{i}{N} + \frac{i}{N}(N-i)\frac{(N+i)u + 3v}{6}\beta, \quad (20)$$

which has been derived for a variety of evolutionary processes before [8, 10, 12, 35, 36, 42].

Next, we address the weak selection approximation of the fixation times. The expectation value of the unconditional fixation time of a single A mutant in a population of B is in general given by the exact equation (14). With the transition and fixation probabilities of the Fermi process, the unconditional fixation time of absorption at any boundary simplifies to

$$t_1 = \phi_1 \sum_{k=1}^{N-1} \sum_{l=1}^k \frac{N^2}{l(N-l)} (1 + e^{-\beta(u l + v)}) \times \exp \left\{ -\beta \sum_{m=l+1}^k \Delta\pi(m) \right\}. \quad (21)$$

The weak selection approximation takes the remarkably simple form (see Appendix B for details)

$$t_1 \approx 2NH_{N-1} + vN(N-1-H_{N-1})\beta, \quad (22)$$

with v given in (6). Thus, t_1 depends only on the constant term of the payoff difference. For large N , this yields $v \approx b - d$. That is, for large populations under weak selection the linear correction of the average fixation time only depends on the advantage (or disadvantage) of the A mutants in the resident population. For $b > d$, invasion of A mutants is likely and slows down the time until the population is homogeneous again. For $d > b$, it is difficult for A to invade a B population and extinction of the mutants is faster than in the neutral case. Note that the payoff entries a and c have no influence on the unconditional fixation time under weak selection corrections. Since fixation is unlikely for weak selection (the probability of fixation of a single A mutant is approximately N^{-1}), the unconditional fixation time

is dominated by the fixation to B . In this case, it is enough to discuss the invasion of A mutants.

Next, we address the average time to fixation given that the A mutant takes over the population. With the general result (15) the Fermi processes conditional fixation time to all A reads

$$t_1^A = \sum_{k=1}^{N-1} \sum_{l=1}^k \phi_l \frac{N^2}{l(N-l)} (1 + e^{-\beta(u l + v)}) \times \exp \left\{ -\beta \sum_{m=1+l}^k \Delta \pi(m) \right\}. \quad (23)$$

Its linear approximation turns out to be dependent on the payoffs in a very simple way as well,

$$t_1^A \approx 2N(N-1) - uN(N-1) \frac{N^2 + N - 6}{18} \beta, \quad (24)$$

with $u(N-1) = a - b - c + d$. The detailed calculation can be found in Appendix B. Since during the fixation process all payoffs are of importance, it is obvious that they all enter here. For example, when it is easy to invade because few mutants have an advantage ($b > d$), but difficult to reach fixation because mutants are disadvantageous once they are frequent ($c > a$), we have $u < 0$ and the conditional time to fixation is larger than neutral. In the last section, we discuss special classes of games to show that, under weak selection, the conditional mean exit times of fixation (or absorption) do not always follow the intuition based on the payoff matrix (1).

6. Frequency dependent Moran process

In this section we address the generality of the previous findings discussing an alternative evolutionary process. The first model that connects payoffs from a 2×2 game to reproductive fitness using a weak selection approach in finite populations is the frequency dependent Moran process [8, 9]. In this process, an individual is chosen for reproduction with probability proportional to its fitness $f(i)$. The offspring replaces a randomly chosen individual. The average payoffs (2) and (3) are mapped to the fitness such that $f_A(i) = 1 - \beta + \beta \pi_A(i)$ and $f_B(i) = 1 - \beta + \beta \pi_B(i)$, where the selection intensity $\beta > 0$ is so small that $f_A(i) > 0$ and $f_B(i) > 0$. The transition

probabilities of the standard Moran process read

$$T_i^+ = \frac{if_A(i)}{if_A(i) + (N-i)f_B(i)} \frac{N-i}{N}, \quad (25)$$

$$T_i^- = \frac{(N-i)f_B(i)}{if_A(i) + (N-i)f_B(i)} \frac{i}{N}. \quad (26)$$

Although these transition probabilities are different from those of the Fermi process, they also yield $\gamma_i \approx 1 - \Delta\pi(i)$ and $\prod_{m=l+1}^k \gamma_m \approx 1 - \beta \sum_{m=l+1}^k \Delta\pi(m)$ for weak selection, $\beta \ll 1$. Thus, the weak selection approximations of the fixation probabilities ϕ_l of the Moran process and the Fermi process are identical, see equation (20). But the weak selection approximation of the transition probabilities are not identical, which leads consequently to different mean exit times. Nevertheless, the results have the same, remarkably simple connection to the payoff matrix (1). The mean exit times or fixation times of the frequency dependent Moran process are

$$t_1 \approx NH_{N-1} + v \frac{N}{2} (N+1 - 2H_N) \beta, \quad (27)$$

$$t_1^A \approx N(N-1) - u \frac{N^2(N^2 - 3N + 2)}{36} \beta. \quad (28)$$

Qualitatively, the dependence on the payoff matrix via u and v is the same as for the Fermi process. Their calculation is analogous to the findings of the previous section, details can be found in Appendix B. Note that, comparing with the Fermi process, there is a factor of 2 missing in the neutral terms. However, this can be avoided by rescaling the transition probabilities, without changing the properties of the different processes.

7. Discussion

Finally, let us discuss the implications of our results for general 2×2 games. While we concentrate on the Fermi process here, the discussion is equally valid for the frequency dependent Moran process. An important question is whether the linear correction for weak selection is compatible with the general features of the game and the known asymptotic behavior for large N of the mean exit or fixation times derived by Antal and Scheuring [26]. Clearly, this depends on the payoff matrix of the 2×2

game,

$$\begin{array}{c} A \quad B \\ A \left(\begin{array}{cc} a & b \\ c & d \end{array} \right), \\ B \end{array} \quad (29)$$

as the payoffs enter the first exit times of absorption linearly. To analyze the difference to the neutral case we consider the rescaled average times of fixation, $\tau_1(\beta) = t_1(\beta)/t_1(0)$ and $\tau_1^A(\beta) = t_1^A(\beta)/t_1^A(0)$. The rescaled unconditional fixation time reads

$$\tau_1 \approx 1 + \frac{1}{2} \frac{N(b-d) - a + d}{N-1} \left(\frac{N-1}{H_{N-1}} - 1 \right) \beta. \quad (30)$$

Accordingly, the rescaled conditional fixation time for absorption at all A is

$$\tau_1^A \approx 1 - \frac{a-b-c+d}{N-1} \frac{N^2 + N - 6}{36} \beta. \quad (31)$$

Note that for population sizes $N > 2$ and sufficiently small β , we always have $t_1(0) < t_1^A(0)$. In other words, the average time until the A individual has reached fixation or gone extinct is smaller than the conditional average time until the A individual has reached fixation. For $\beta \rightarrow \infty$, the process follows deterministically the intensity of selection and thus both fixation times may coincide, $t_1(\beta \rightarrow \infty) \approx t_1^A(\beta \rightarrow \infty)$. This ordering of the fixation times is blurred by our rescaling, as we focus only on the change relative to the neutral case.

In the following, we discuss these two expressions for the three generic types of 2×2 games, namely dominance of A ($a > c$ and $b > d$), coexistence of A and B ($a < b$ and $c > d$) and a coordination game ($a > c$ and $b < d$).

7.1. Dominance of A .

Consider a game where strategy A is always dominant, i.e. it obtains a larger payoff than B , regardless of the fraction of A in the population. This is the case for $a > c$ and $b > d$. One special case is the Prisoner's Dilemma with $b > d > a > c$. The interesting feature of this game is that the social optimum d is not the Nash equilibrium, which is a . For neutral selection, a single A individual goes extinct with probability $1 - N^{-1}$. Thus, the unconditional fixation time τ_1 is dominated by the extinction of A . Since strategy A is favored by selection, increasing the intensity of selection decreases the probability of the extinction of A . Since fixation takes

at least $N - 1$ time steps, τ_1 increases with increasing intensity of selection β . For large N , this is obvious from our equation (30), because in this case the quantity $N(b-d) - a + d$ is positive. However, once extinction of A becomes unlikely, increasing β further will lead to a decrease of τ_1 .

The discussion of the conditional fixation time τ_1^A is not as straightforward, because the sign of $a - b - c + d$ can be positive or negative. The sign of this quantity is also decisive for the evolutionary dynamics in other contexts, see e.g. [43]. When the advantage of an A individual is initially large and decreases with the abundance of A ($a - c > b - d > 0$), then the sign of $a - b - c + d$ is positive and τ_1^A decreases with increasing intensity of selection. But when the advantage of strategy A decreases with the number of A individuals ($b - d > a - c > 0$), then τ_1^A increases with increasing intensity of selection. However, this apparently counterintuitive phenomenon (after all, A dominates B) can only be observed for weak selection. For strong selection, τ_1^A decreases again. These results are compatible with the observation that the conditional fixation time scales as $N \ln N$ for large N [26]. In Figure 2 (a) we show a numerical example for the rescaled average times. We include averages from numerical simulations of the evolutionary process, our linear approximation as well as the exact result that can be obtained from dividing equation (14) by (18) and equation (15) by (19), respectively. The payoff matrix is chosen such that $a + d > b + c$, which means that with increasing intensity of selection τ_1^A decreases and τ_1 increases.

7.2. Coexistence of A and B .

As a second class, we consider games in which B is the best reply to A ($c > a$), but A is the best reply to B ($b > d$). Important examples for such games are the Hawk-Dove game [44] or the Snowdrift game [45]. For infinite populations, the replicator dynamics predicts a stable coexistence of A and B . In finite populations, the system typically fluctuates around that point until eventually, fluctuations lead to absorption in one of the boundaries [46, 47]. Consequently, the conditional fixation times increase exponentially with the population size [26]. Since $a - b - c + d$ is negative, we also have an increase of τ_1^A with the selection intensity for weak selection. Further, $N(b - d) - a + d$ is positive in large populations, such that also τ_1 increases with the selection intensity. Figure 2 (b) shows that the divergence of the exact results is faster than the linear approximation even for weak selection.

7.3. Coordination games.

Finally, let us discuss coordination games in which $a > c$ and $b < d$. In these games, A is the best reply to A and B is the best reply to B . The replicator equation of such systems exhibits a bistability: If the fraction of A individuals is sufficiently high in the beginning, the A individuals will reach fixation. Otherwise, B individuals will take over the system. The stronger the intensity of selection, the less likely it is that a single A individual can take over a B population. Consequently, τ_1 should decrease with β . This also follows from our weak selection approximation: In large populations, $N(b - d) - a + d$ is negative and thus τ_1 decreases with the intensity of selection, see equation (30). Perhaps less intuitive, also τ_1^A decreases with β , which results from $a - b - c + d > 0$, cf. (30). However, this is again consistent with the observation that τ_1^A scales as $N \ln N$ in large populations. Although the fixation probability of a single A decreases with β , if such an event occurs, it is faster than in the neutral case. A numerical example for this behavior is shown in Figure 2 (c).

The numerical examples indicate that the convergence radius of our weak selection expansion is of the order of N^{-1} , which is also known for many systems in population genetics. Although N^{-1} might appear small, this kind of weak selection is the most relevant limit in evolutionary biology, as evolutionary change is typically only connected with small selective differences. We stress that we have made no assumptions on the population size, such that our results are valid for arbitrary N .

Our approach shows under which circumstances the general features of the game are reflected in the fixation times under weak selection. Although the weak selection expansion of the mean exit or fixation times is technically rather tedious, the resulting asymptotic behavior shows remarkable simplicity.

Acknowledgment

Financial support by the Emmy-Noether program of the DFG is gratefully acknowledged.

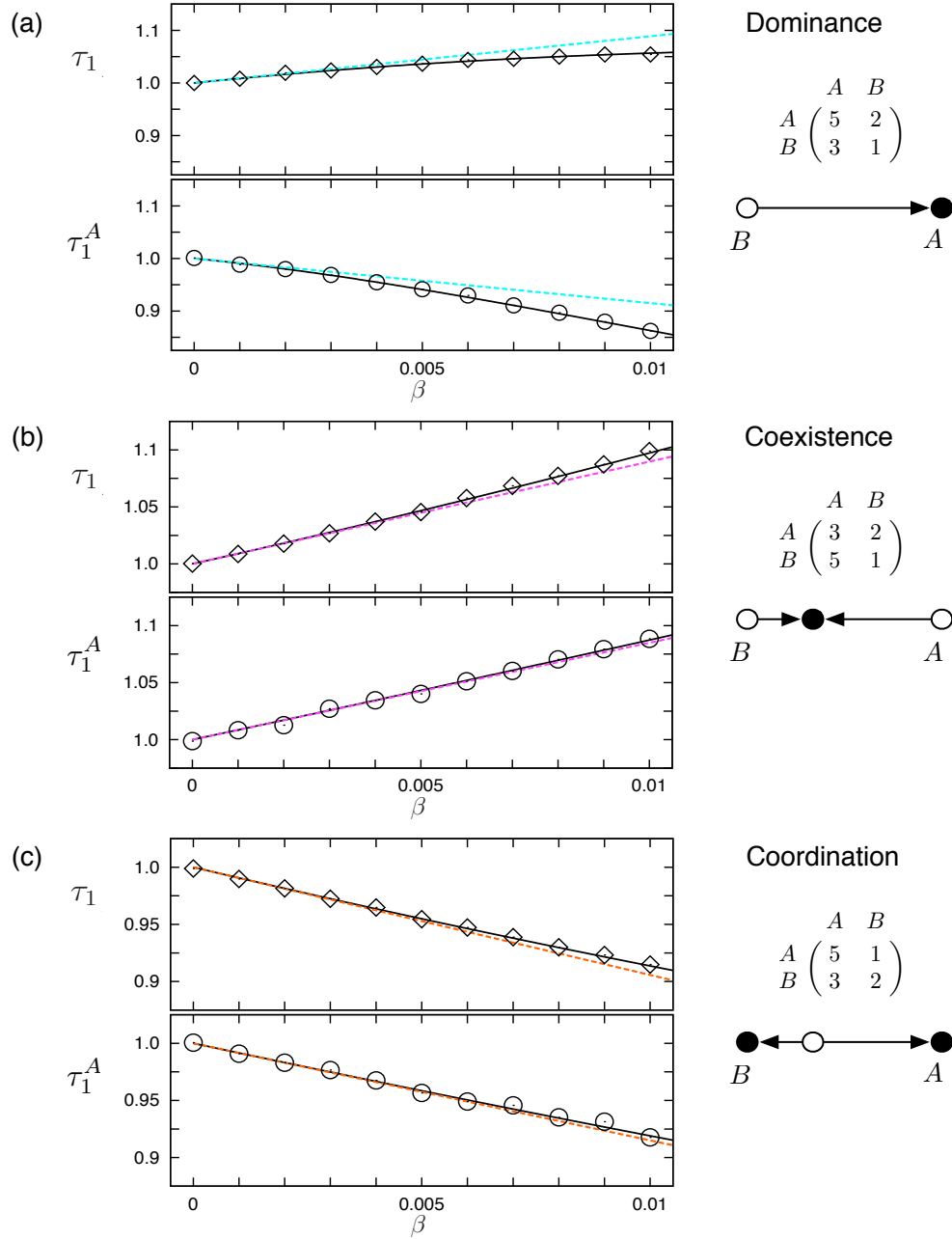


Figure 2: Expectation values of the rescaled fixation times starting with a single A mutant in a population of B as a function of the selection intensity β . Full lines show the normalized exact solution originating from the exact results (21) and (23). Colored dashed lines are the linear approximations (30) and (31). Symbols show the results from simulations based on 10^7 realizations, which agree nicely with the exact results. Diamonds are for the unconditional averages, circles are for the conditional averages. On the right hand side, we show the payoff matrices of the three games and illustrate the direction of selection in these games. (a) In a game with dominance of strategy A , the unconditional fixation time increases with the intensity of selection, but the conditional fixation time decreases. (b) For games with stable coexistence, both fixation times increase with the intensity of selection. (c) For coordination games, the two fixation times become shorter when the intensity of selection is increased. In all examples, the population size is $N = 100$.

Appendix A. Finite double sums

Here, we collect some helpful calculations for double sums as they appear in the mean exit times. An important observation is

$$\sum_{k=i}^{N-1} \sum_{l=1}^k \frac{f_l}{N-l} = (N-i) \sum_{l=1}^{i-1} \frac{f_l}{N-l} + \sum_{l=i}^{N-1} f_l, \quad (\text{A.1})$$

for any function $f_l < \infty$ and $l = 1, \dots, N-1$. This can be seen by writing the left hand side term by term, i.e.

$$\begin{aligned} \sum_{k=i}^{N-1} \sum_{l=1}^k \frac{f_l}{N-l} &= \frac{f_1}{N-1} + \dots + \frac{f_i}{N-i} \\ &\quad + \frac{f_1}{N-1} + \dots + \frac{f_i}{N-i} + \frac{f_{i+1}}{N-(i+1)} \\ &\quad + \dots \\ &\quad + \frac{f_1}{N-1} + \dots + \frac{f_i}{N-i} + \dots + \frac{f_{N-1}}{N-(N-1)} \\ &= (N-i) \sum_{l=1}^i \frac{f_l}{N-l} \\ &\quad + (N-i-1) \frac{f_{i+1}}{N-(i+1)} + \dots + f_{N-1} \\ &= (N-i) \sum_{l=1}^{i-1} \frac{f_l}{N-l} + \sum_{l=i}^{N-1} f_l. \end{aligned} \quad (\text{A.2})$$

For the case $i = 1$ the result is especially simple, since the first sum of the right hand side of equation (A.1) vanishes. This case is of special interest for the computation of t_1^A under neutral selection with $f_l = 1$ and for t_1 with $f_l = 1/l$.

Another finding for double sums with $M \in \mathbb{N}$ and two bounded functions f_k and g_l is

$$\sum_{k=1}^M \sum_{l=1}^k f_k g_l = \sum_{l=1}^M g_l \sum_{k=l}^M f_k. \quad (\text{A.3})$$

This becomes clear by resorting the terms again,

$$\sum_{k=1}^M \sum_{l=1}^k f_k g_l = f_1 g_1 + f_2 (g_1 + g_2) + \dots + f_M (g_1 + g_2 + \dots + g_M)$$

$$\begin{aligned}
&= g_1(f_1 + \dots + f_M) + g_2(f_2 + \dots + f_M) + \dots + g_M f_M \\
&= \sum_{l=1}^M g_l \sum_{k=l}^M f_k.
\end{aligned} \tag{A.4}$$

Appendix B. Fixation times under weak selection

Here, we calculate the linear corrections of the mean exit times t_1 and t_1^A for the Fermi process in detail, compare equations (21) and (23). We aim at finding these times for weak selection, e.g.

$$t_1 \approx [t_1]_{\beta=0} + \beta \left[\frac{\partial}{\partial \beta} t_1 \right]_{\beta=0}. \tag{B.1}$$

The first term follows directly from the calculation in Appendix A, see equation (18). Our goal here is to compute the linear term $\left[\frac{\partial}{\partial \beta} t_1 \right]_{\beta=0}$.

$$\begin{aligned}
\left[\frac{\partial}{\partial \beta} t_1 \right]_{\beta=0} &= \sum_{k=1}^{N-1} \sum_{l=1}^k \left[\frac{1}{T_l^+} \frac{\partial \phi_1}{\partial \beta} + \phi_1 \frac{\partial}{\partial \beta} \frac{1}{T_l^+} \right]_{\beta=0} \\
&\quad - \sum_{k=1}^{N-1} \sum_{l=1}^k \left[\frac{\phi_1}{T_l^+} \sum_{m=l+1}^k \Delta \pi(m) \right]_{\beta=0},
\end{aligned} \tag{B.2}$$

where we applied $\left[\prod_{m=l+1}^k \gamma_m \right]_{\beta=0} = 1$ and $\left[\frac{\partial}{\partial \beta} \prod_{m=l+1}^k \gamma_m \right]_{\beta=0} = - \sum_{m=l+1}^k \Delta \pi(m)$. For the fixation probability under weak selection and with $\Delta \pi(l) = ul + v$, we have

$$\left[\frac{\partial \phi_l}{\partial \beta} \right]_{\beta=0} = \frac{l}{N} (N-l) \frac{(N+l)u + 3v}{6}. \tag{B.3}$$

The weak selection approximation of the inverse of the transition probability T_l^+ , compare equation (8), yields

$$\left[\frac{\partial}{\partial \beta} \frac{1}{T_l^+} \right]_{\beta=0} = - \frac{N^2}{l(N-l)} (ul + v). \tag{B.4}$$

This leads to

$$\begin{aligned}
\left[\frac{\partial}{\partial \beta} t_1 \right]_{\beta=0} &= \sum_{k=1}^{N-1} \sum_{l=1}^k \frac{(N-1)((N+1)u + 3v)}{6N} \frac{2N^2}{l(N-l)} \\
&\quad - \sum_{k=1}^{N-1} \sum_{l=1}^k \frac{1}{N} \frac{N^2}{l(N-l)} (ul + v)
\end{aligned}$$

$$- \sum_{k=1}^{N-1} \sum_{l=1}^k \frac{1}{N} \frac{2N^2}{l(N-l)} \sum_{m=l+1}^k (u m + v) \quad (\text{B.5})$$

While the first two double sums can be solved with the help of Appendix A, the third term is more complicated. For this more tedious calculation, we refer to Appendix C. Eventually, the solution of the double and triple sums leads to

$$\begin{aligned} \left[\frac{\partial}{\partial \beta} t_1 \right]_{\beta=0} &= N(N-1) \frac{(N+1)u + 3v}{3} H_{N-1} \\ &\quad - N(N-1)u - N H_{N-1} v \\ &\quad - N(N-1) \left(\frac{((N+1)u + 3v)}{3} H_{N-1} - u - v \right) \\ &= v N(N-1 - H_{N-1}), \end{aligned} \quad (\text{B.6})$$

where the last step is elementary. Combining this with equation (18) leads finally to the unconditional mean exit time under weak selection, equation (22).

For the conditional fixation time t_1^A , the linear term $\left[\frac{\partial}{\partial \beta} t_1^A \right]_{\beta=0}$ reads

$$\begin{aligned} \left[\frac{\partial}{\partial \beta} t_1^A \right]_{\beta=0} &= \sum_{k=1}^{N-1} \sum_{l=1}^k \left[\frac{1}{T_l^+} \frac{\partial \phi_l}{\partial \beta} + \phi_l \frac{\partial}{\partial \beta} \frac{1}{T_l^+} \right]_{\beta=0} \\ &\quad - \sum_{k=1}^{N-1} \sum_{l=1}^k \left[\frac{\phi_l}{T_l^+} \sum_{m=l+1}^k \Delta \pi(m) \right]_{\beta=0}. \end{aligned} \quad (\text{B.7})$$

The only difference compared to the unconditional fixation time, equation (B.2), is the fixation probability ϕ_l instead of ϕ_1 . The linear term of the weak selection expansion of ϕ_l is given in equation (B.3). This yields

$$\begin{aligned} \left[\frac{\partial}{\partial \beta} t_1^A \right]_{\beta=0} &= \sum_{k=1}^{N-1} \sum_{l=1}^k \frac{l(N-l)((N+l)u + 3v)}{6N} \frac{2N^2}{l(N-l)} \\ &\quad - \sum_{k=1}^{N-1} \sum_{l=1}^k \frac{l}{N} \frac{N^2}{l(N-l)} (u l + v) \\ &\quad - \sum_{k=1}^{N-1} \sum_{l=1}^k \frac{l}{N} \frac{2N^2}{l(N-l)} \sum_{m=l+1}^k (u m + v). \end{aligned} \quad (\text{B.8})$$

Again, the first two double sums can be solved using the results from Appendix A. The third term follows from a calculation which is similar to Appendix C, but

simpler. This last term reduces to

$$\begin{aligned} & \sum_{k=1}^{N-1} \sum_{l=1}^k \frac{l}{N} \frac{2N^2}{l(N-l)} \sum_{m=l+1}^k (u m + v) \\ &= N \frac{(N-2)(N-1)}{18} ((5N+3)u + 9v). \end{aligned} \quad (\text{B.9})$$

Finally, combining the three terms again results in

$$\begin{aligned} \left[\frac{\partial}{\partial \beta} t_1^A \right]_{\beta=0} &= \frac{N^2(N-1)}{18} ((4N+1)u + 9v) \\ &\quad - \frac{N(N-1)}{2} (Nu + 2v) \\ &\quad - (N-2) \frac{N(N-1)}{18} ((5N+3)u + 9v) \\ &= -u N(N-1) \frac{N^2 + N - 6}{18}. \end{aligned} \quad (\text{B.10})$$

In combination with equation (19), this results in the conditional mean exit time under weak selection, equation (24).

For completeness, we briefly repeat this calculation for the mean exit times of the frequency dependent Moran process. With the transition probabilities (25) and (26), the fixation probabilities under weak selection are identical to those of the Fermi process, see equation (20). However, the inverse transition probability is different in the weak selection regime, i.e. the linear correction is

$$\left[\frac{\partial}{\partial \beta} \frac{1}{T_l^+} \right]_{\beta=0} = -\frac{N}{l} \Delta \pi(l) = -N \frac{ul + v}{l}. \quad (\text{B.11})$$

Hence, for the unconditional mean exit time we have the same starting equation (B.2). But with equation (B.11) this gives

$$\begin{aligned} \left[\frac{\partial}{\partial \beta} t_1 \right]_{\beta=0} &= \sum_{k=1}^{N-1} \sum_{l=1}^k \frac{(N-1)((N+1)u + 3v)}{6N} \frac{N^2}{l(N-l)} \\ &\quad - \sum_{k=1}^{N-1} \sum_{l=1}^k \frac{1}{N} N \frac{ul + v}{l} \\ &\quad - \sum_{k=1}^{N-1} \sum_{l=1}^k \frac{1}{N} \frac{N^2}{l(N-l)} \sum_{m=1+l}^k (u m + v), \end{aligned} \quad (\text{B.12})$$

which differs from equation (B.5) only in the second double sum. With the previous findings for the Fermi processes times the required calculation is straightforward and results in

$$\left[\frac{\partial}{\partial \beta} t_1 \right]_{\beta=0} = v \frac{N}{2} (N + 1 - 2H_N). \quad (\text{B.13})$$

That is, this linear correction has a different dependence on the system size N .

For the conditional mean exit time the situation is similar. In difference to equation (B.8), the linear correction reads

$$\begin{aligned} \left[\frac{\partial}{\partial \beta} t_1^A \right]_{\beta=0} &= \sum_{k=1}^{N-1} \sum_{l=1}^k \frac{l(N-l)((N+l)u + 3v)}{6N} \frac{N^2}{l(N-l)} \\ &\quad - \sum_{k=1}^{N-1} \sum_{l=1}^k \frac{l}{N} N \frac{ul + v}{l} \\ &\quad - \sum_{k=1}^{N-1} \sum_{l=1}^k \frac{l}{N} \frac{2N^2}{l(N-l)} \sum_{m=l+1}^k (um + v). \end{aligned} \quad (\text{B.14})$$

This leads to

$$\left[\frac{\partial}{\partial \beta} t_1^A \right]_{\beta=0} = -u \frac{N^2}{36} (N^2 - 3N + 2), \quad (\text{B.15})$$

for the linear correction of the conditional mean exit times of the frequency dependent Moran process.

Appendix C. Finite triple sum

Here, we calculate the triple sum from Appendix B, that require some additional steps. Our goal is to solve

$$\sigma = \sum_{k=1}^{N-1} \sum_{l=1}^k \frac{1}{l(N-l)} \sum_{m=1+l}^k \Delta\pi(m). \quad (\text{C.1})$$

For the sum over payoff differences, we have

$$\sum_{m=1+l}^k \Delta\pi(m) = \sum_{m=1+l}^k (um + v) = f_k - f_l, \quad (\text{C.2})$$

where we introduced the function

$$f_m = m(m+1) \frac{u}{2} + m v, \quad (\text{C.3})$$

which is valid for any integer m . Using partial fraction expansion, $\frac{N}{l(N-l)} = \frac{1}{l} + \frac{1}{N-l}$, we obtain

$$\begin{aligned}\sigma &= \sum_{k=1}^{N-1} \sum_{l=1}^k \frac{f_k}{l(N-l)} - \sum_{k=1}^{N-1} \sum_{l=1}^k \frac{f_l}{l(N-l)} \\ &= \underbrace{\frac{1}{N} \sum_{k=1}^{N-1} \sum_{l=1}^k \frac{f_k}{l}}_{K_1} + \underbrace{\frac{1}{N} \sum_{k=1}^{N-1} \sum_{l=1}^k \frac{f_k}{N-l}}_{K_2} \\ &\quad - \underbrace{\frac{1}{N} \sum_{k=1}^{N-1} \sum_{l=1}^k \frac{f_l}{l}}_{K_3} - \underbrace{\frac{1}{N} \sum_{k=1}^{N-1} \sum_{l=1}^k \frac{f_l}{N-l}}_{K_4}.\end{aligned}\tag{C.4}$$

We solve each part separately, starting with the last one. For K_4 , we obtain with equation (A.1) from Appendix A

$$K_4 = \sum_{k=1}^{N-1} \sum_{l=1}^k \frac{f_l}{N-l} = \sum_{k=1}^{N-1} f_k = \frac{N-1}{6} N((N+1)u + 3v).\tag{C.5}$$

The second last term, K_3 , is a sum over a linear function and can be treated with any table of elementary sums, e. g. [48],

$$\begin{aligned}K_3 &= \sum_{k=1}^{N-1} \sum_{l=1}^k \frac{f_l}{l} = \sum_{k=1}^{N-1} \sum_{l=1}^k ((l+1)\frac{u}{2} + v) \\ &= \frac{N-1}{12} N((N+4)u + 6v).\end{aligned}\tag{C.6}$$

The remaining two terms require more effort. Both terms, K_1 and K_2 have the same structure regarding functions of k and l . Using equation (A.3), we have

$$K_2 = \sum_{k=1}^{N-1} \sum_{l=1}^k \frac{f_k}{N-l} = \sum_{l=1}^{N-1} \frac{1}{N-l} \sum_{k=l}^{N-1} f_k,\tag{C.7}$$

and

$$K_1 = \sum_{k=1}^{N-1} \sum_{l=1}^k \frac{f_k}{l} = \sum_{l=1}^{N-1} \frac{1}{l} \sum_{k=l}^{N-1} f_k.\tag{C.8}$$

Hence, we first have to compute the sum $\sum_{k=l}^{N-1} f_k$, which reduces to the solution of elementary sums,

$$\sum_{k=l}^{N-1} f_k = \sum_{k=l}^{N-1} k((k+1)\frac{u}{2} + v) = \frac{u}{2} \sum_{k=l}^{N-1} k^2 + \left(\frac{u}{2} + v\right) \sum_{k=l}^{N-1} k$$

$$\begin{aligned}
&= \frac{N-l}{6} (N^2 + Nl + l^2 - 1)u + 3(N+l-1)v \\
&= \frac{N-l}{6} (N^2 + Nl + l^2 - 1)u + \frac{N-l}{2} (N+l-1)v. \quad (C.9)
\end{aligned}$$

Thus, solving equations (C.7) and (C.8) simplifies to solving the elementary sums $\sum_{l=1}^{N-1} l^s$ with $s = 0, 1, 2$, compare [48]. With this, we have

$$\begin{aligned}
K_2 &= \frac{N-1}{6} \sum_{l=1}^{N-1} ((N+1)u + 3v) + \frac{Nu + 3v}{6} \sum_{l=1}^{N-1} l + \frac{u}{6} \sum_{l=1}^{N-1} l^2 \\
&= \frac{N-1}{36} ((11N^2 - N - 6)u + 9(3N - 2)v). \quad (C.10)
\end{aligned}$$

For K_1 , we obtain

$$\begin{aligned}
K_1 &= \frac{1}{6} \sum_{l=1}^{N-1} \frac{N(N-1)(N+1)u + 3N(N-1)v}{l} \\
&\quad - \frac{1}{6} \sum_{l=1}^{N-1} ((l^2 - 1)u - 3(l-1)v) \\
&= \frac{N(N-1)}{6} ((N+1)u + 3v) H_{N-1} \\
&\quad - \frac{N-1}{36} (N-2)((2N+3)u + 9v).
\end{aligned}$$

Summing up the terms, $\sigma = (K_1 + K_2 - K_3 - K_4)/N$, finally yields the result

$$\sigma = \frac{N-1}{6} (((N+1)u + 3v)H_{N-1} - 3(u+v)). \quad (C.11)$$

Again, $H_n = \sum_{l=1}^n 1/l$ are the harmonic numbers. In equation (B.9), the reasoning is very similar, but only terms of the structure of K_2 and K_4 appear.

References

- [1] G. Szabó and G. Fáth. Evolutionary games on graphs. *Physics Reports*, 446:97–216, 2007.
- [2] J. Berg and A. Engel. Matrix games, mixed strategies, and statistical mechanics. *Phys. Rev. Lett.*, 81:4999–5002, 1998.
- [3] G. Szabó and C. Hauert. Phase transitions and volunteering in spatial public goods games. *Phys. Rev. Lett.*, 89:118101, 2002.
- [4] F. C. Santos and J. M. Pacheco. Scale-free networks provide a unifying framework for the emergence of cooperation. *Phys. Rev. Lett.*, 95:098104, 2005.
- [5] C. Hauert and G. Szabó. Game theory and physics. *Am. Journal of Physics*, 73:405–414, 2005.

- [6] P. D. Taylor and L. Jonker. Evolutionary stable strategies and game dynamics. *Math. Biosci.*, 40:145–156, 1978.
- [7] J. Hofbauer and K. Sigmund. *Evolutionary Games and Population Dynamics*. Cambridge University Press, Cambridge, 1998.
- [8] M. A. Nowak, A. Sasaki, C. Taylor, and D. Fudenberg. Emergence of cooperation and evolutionary stability in finite populations. *Nature*, 428:646–650, 2004.
- [9] C. Taylor, D. Fudenberg, A. Sasaki, and M. A. Nowak. Evolutionary game dynamics in finite populations. *Bull. Math. Biol.*, 66:1621–1644, 2004.
- [10] A. Traulsen, J. C. Claussen, and C. Hauert. Coevolutionary dynamics: From finite to infinite populations. *Phys. Rev. Lett.*, 95:238701, 2005.
- [11] G. Wild and P.D. Taylor. Fitness and evolutionary stability in game theoretic models of finite populations. *Proc. Roy. Soc. Lond. B*, 271:2345–2349, 2004.
- [12] L. A. Imhof, D. Fudenberg, and M. A. Nowak. Evolutionary cycles of cooperation and defection. *Proc. Natl. Acad. Sci. USA*, 102:10797–10800, 2005.
- [13] A. Traulsen, J. M. Pacheco, and L. A. Imhof. Stochasticity and evolutionary stability. *Phys. Rev. E*, 74:021905, 2006.
- [14] D. Fudenberg, M. A. Nowak, C. Taylor, and L.A. Imhof. Evolutionary game dynamics in finite populations with strong selection and weak mutation. *Theor. Pop. Biol.*, 70:352–363, 2006.
- [15] T. Reichenbach, M. Mobilia, and E. Frey. Coexistence versus extinction in the stochastic cyclic Lotka-Volterra model. *Phys. Rev. E*, 74:051907, 2006.
- [16] M. Perc. Coherence resonance in a spatial prisoner’s dilemma game. *New J. Physics*, 8:22–33, 2006.
- [17] M. Perc and M. Marhl. Evolutionary and dynamical coherence resonances in the pair approximated prisoner’s dilemma game. *New J. Physics*, 8:142, 2006.
- [18] M. Perc and A. Szolnoki. Noise-guided evolution within cyclical interactions. *New J. Physics*, 9:267, 2007.
- [19] J. C. Claussen. Drift reversal in asymmetric coevolutionary conflicts: influence of microscopic processes and population size. *European Physical Journal B*, 60:391–399, 2007.
- [20] J. Cremer, T. Reichenbach, and E. Frey. Anomalous finite-size effects in the battle of the sexes. *European Physical Journal B*, 63(3), 2008.
- [21] A. Szolnoki and M. Perc. Coevolution of teaching activity promotes cooperation. *New J. Physics*, 10:043036, 2008.
- [22] J. C. Claussen and A. Traulsen. Cyclic dominance and biodiversity in well-mixed populations. *Phys. Rev. Lett.*, 100:058104, 2008.
- [23] T. Ohta. Near-neutrality in evolution of genes and gene regulation. *Proc. Natl. Acad. Sci. USA*, 99:16134–16137, 2002.
- [24] H. Ohtsuki, C. Hauert, E. Lieberman, and M. A. Nowak. A simple rule for the evolution of cooperation on graphs. *Nature*, 441:502–505, 2006.
- [25] H. Ohtsuki, M. A. Nowak, and J. M. Pacheco. Breaking the symmetry between interaction and replacement in evolutionary dynamics on graphs. *Phys. Rev. Lett.*, 98:108106, 2007.
- [26] T. Antal and I. Scheuring. Fixation of strategies for an evolutionary game in finite populations. *Bull. Math. Biol.*, 36(12):1923–1944, 2006.

- [27] N. G. van Kampen. *Stochastic Processes in Physics and Chemistry*. Elsevier, Amsterdam, 2 edition, 1997.
- [28] S. Redner. *A Guide to First-Passage Processes*. Cambridge University Press, 2001.
- [29] C. W. Gardiner. *Handbook of Stochastic Methods*. Springer New York, third edition, 2004.
- [30] Lawrence E. Blume. The statistical mechanics of strategic interaction. *Games and Economic Behavior*, 4:387–424, 1993.
- [31] G. Szabó and C. Tóke. Evolutionary Prisoner’s Dilemma game on a square lattice. *Phys. Rev. E*, 58:69, 1998.
- [32] J. M. Pacheco, A. Traulsen, and M. A. Nowak. Co-evolution of strategy and structure in complex networks with dynamical linking. *Phys. Rev. Lett.*, 97:258103, 2006.
- [33] A. Traulsen, J. M. Pacheco, and M. A. Nowak. Pairwise comparison and selection temperature in evolutionary game dynamics. *J. Theor. Biol.*, 246:522–529, 2007.
- [34] S. Karlin and H. M. A. Taylor. *A first course in stochastic processes*. Academic, London, 2nd edition edition, 1975.
- [35] A. Traulsen and C. Hauert. Stochastic evolutionary game dynamics. In H.-G. Schuster, editor, *Reviews of nonlinear dynamics and complexity*. Wiley New York, 2009, arXiv:0811.3538.
- [36] M. A. Nowak. *Evolutionary Dynamics*. Harvard University Press, Cambridge, MA, 2006.
- [37] A. Traulsen, M. A. Nowak, and J. M. Pacheco. Stochastic dynamics of invasion and fixation. *Phys. Rev. E*, 74:11909, 2006.
- [38] M. E. Fisher. Diffusion from an entrance to an exit. *IBM J. Res. Dev.*, 32:76–81, 1988.
- [39] N.S. Goel and N. Richter-Dyn. *Stochastic Models in Biology*. Academic Press, New York, 1974.
- [40] D. Dingli, A. Traulsen, and J. M. Pacheco. Stochastic dynamics of hematopoietic tumor stem cells. *Cell Cycle*, 6:e2–e6, 2007.
- [41] M. Kimura. Evolutionary rate at the molecular level. *Nature*, 217:624–626, 1968.
- [42] S. Lessard and V. Ladret. The probability of fixation of a single mutant in an exchangeable selection model. *J. Math. Biol.*, 54:721–744, 2007.
- [43] C. Taylor and M. A. Nowak. Transforming the dilemma. *Evolution*, 61:2281–2292, 2007.
- [44] J. Maynard Smith and G. R. Price. The logic of animal conflict. *Nature*, 246:15–18, 1973.
- [45] M. Doebeli and C. Hauert. Models of cooperation based on the prisoner’s dilemma and the snowdrift game. *Ecology Letters*, 8:748–766, 2005.
- [46] J. C. Claussen and A. Traulsen. Non-Gaussian fluctuations arising from finite populations: Exact results for the evolutionary Moran process. *Phys. Rev. E*, 71:025101(R), 2005.
- [47] F. A. C. C. Chalub and M. O. Souza. Discrete versus continuous models in evolutionary dynamics: From simple to simpler – and even simpler – models. *Math. and Comp. Modelling*, 47:743–754, 2008.
- [48] R. L. Graham, D. E. Knuth, and O. Patashnik. *Concrete Mathematics*. Addison-Wesley, second edition, 1994.